

# On the problem of uniqueness for the steady Navier–Stokes equation in a cascade of profiles

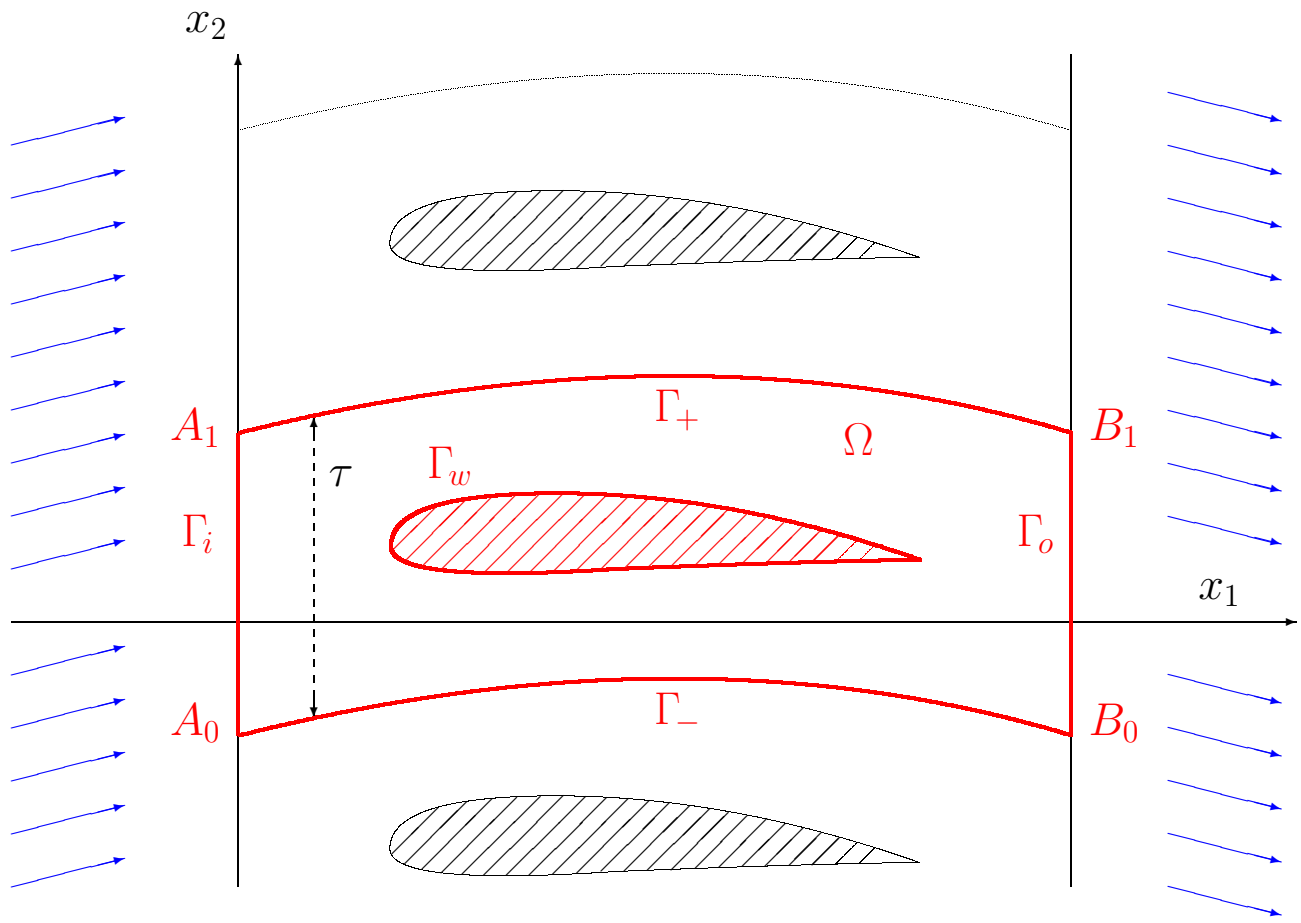
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- Geometry of the problem
  - The problem with linear boundary conditions
  - The problem with a nonlinear boundary condition on the outflow

We consider a flow through the cascade of profiles. We assume that the fluid is **viscous, incompressible, Newtonian** and the flow is **steady** and **2D**.



# 1. The problem with linear boundary conditions

## 1.1. The equations of motion

The conservation of momentum is expressed by the **Navier–Stokes equation**

$$\omega(\mathbf{u}) \mathbf{u}^\perp = -\nabla q + \nu (-\partial_2, \partial_1) \omega(\mathbf{u}) + \mathbf{f} \quad (1)$$

where  $q = p + \frac{1}{2} (u_1^2 + u_2^2)$  is the so-called **Bernoulli pressure**,  $\omega(\mathbf{u}) = \partial_1 u_2 - \partial_2 u_1$  is the **vorticity**, and  $\mathbf{u}^\perp = (-u_2, u_1)$ .

$\mathbf{u} = (u_1, u_2)$	...	<b>velocity</b> ,
$p$	...	<b>pressure</b> ,
$\mathbf{f} = (f_1, f_2)$	...	<b>specific volume force</b> ,
$\nu > 0$	...	<b>kinematic viscosity</b> .

The condition of incompressibility is expressed by **the equation of continuity**

$$\operatorname{div} \mathbf{u} = 0.$$

(2)

It also expresses the conservation of mass.

## 1.2 Boundary conditions

The **inhomogeneous Dirichlet condition on the inflow**:

$$\boxed{\mathbf{u} = \mathbf{g}} \quad \text{on } \Gamma_i \tag{3}$$

where  $\mathbf{g}$  is a given velocity on  $\Gamma_i$ .

The **conditions of periodicity on  $\Gamma_-$  and  $\Gamma_+$** :

$$\boxed{\begin{aligned} \mathbf{u}(x_1, x_2 + \tau) &= \mathbf{u}(x_1, x_2) \\ \frac{\partial \mathbf{u}}{\partial \mathbf{n}}(x_1, x_2 + \tau) &= -\frac{\partial \mathbf{u}}{\partial \mathbf{n}}(x_1, x_2) \\ q(x_1, x_2 + \tau) &= q(x_1, x_2) \end{aligned}} \tag{4}$$

$$\text{for } (x_1, x_2) \in \Gamma_- \tag{5}$$

$$\tag{6}$$

The **homogeneous Dirichlet condition** on the profile:

$$\mathbf{u} = \mathbf{0} \quad \text{on } \Gamma_w \quad (7)$$

The **linear "do-nothing" condition** on the outflow  $\Gamma_o$ :

$$q = h_1 \quad -\nu \omega(\mathbf{u}) = h_2 \quad (8)$$

### 1.3 Function spaces and the weak formulation

$H^1(\Omega)$  (respectively  $H^1(\Omega)^2$ ) is the Sobolev space of scalar (respectively vector) functions, defined a.e. in  $\Omega$ , with the norm  $\|\cdot\|_1$ .

$H^s(\Gamma_i)^2$  (for  $0 < s < 1$ ) is the Sobolev–Slobodetski space of vector functions, defined a.e. in  $\Gamma_i$ , with the norm  $\|\cdot\|_{s;\Gamma_i}$ .

$$\begin{aligned} V = \{ & \mathbf{v} \in H^1(\Omega)^2; \quad \mathbf{v} = \mathbf{0} \text{ s.v. v } \Gamma_i \cup \Gamma_w, \\ & \mathbf{v}(x_1, x_2 + \tau) = \mathbf{v}(x_1, x_2) \text{ for a.a. } (x_1, x_2) \in \Gamma_-, \\ & \operatorname{div} \mathbf{v} = 0 \text{ a.e. in } \Omega \}. \end{aligned}$$

(Conditions on  $\Gamma_i$ ,  $\Gamma_w$  a  $\Gamma_-$  are fulfilled in the sense of traces.)

The norm in  $V$  is

$$\|\|\mathbf{v}\|\| := \|\nabla \mathbf{v}\|_{L^2(\Omega)^4}.$$

## Formal derivation of the weak formulation:

In order to derive formally the weak formulation of the problem, we multiply equation (1) by an arbitrary test function  $\mathbf{v} = (v_1, v_2) \in V$ , integrate in  $\Omega$  and use Green's theorem. We obtain

$$a(\mathbf{u}, \mathbf{v}) = (\mathbf{f}, \mathbf{v})_{L^2(\Omega)^2} + b(\mathbf{h}, \mathbf{v}) \quad \forall \mathbf{v} \in V,$$

where

$$a(\mathbf{u}, \mathbf{v}) \quad := \quad a_1(\mathbf{u}, \mathbf{v}) + a_2(\mathbf{u}, \mathbf{u}, \mathbf{v}),$$

$$a_1(\mathbf{u}, \mathbf{v}) \quad := \quad \nu (\omega(\mathbf{u}), \omega(\mathbf{v}))_{L^2(\Omega)},$$

$$a_2(\mathbf{u}, \mathbf{v}, \mathbf{w}) \quad := \quad \int_{\Omega} \omega(\mathbf{u}) \mathbf{v}^{\perp} \cdot \mathbf{w} \, d\mathbf{x},$$

$$b(\mathbf{h}, \mathbf{v}) \quad := \quad - \int_{\Gamma_o} \mathbf{h} \cdot \mathbf{v} \, dS.$$

The term implying the form of the condition on  $\Gamma_o$ :  $\int_{\Gamma_o} [\nu \omega(\mathbf{u}) v_2 - q v_1] \, dS$



The weak formulation of the problem (1)–(8).

**Definition 1.** Let  $\mathbf{g} \in H^s(\Gamma_i)^2$  (for some  $s \in (\frac{1}{2}, 1]$ ) satisfy the condition  $\mathbf{g}(A_1) = \mathbf{g}(A_0)$ . Let  $\mathbf{f} \in L^2(\Omega)^2$  and  $\mathbf{h} \in L^2(\Gamma_o)^2$ . The **weak solution** of the problem (1)–(8) is a vector function  $\mathbf{u} \in H^1(\Omega)^2$  which satisfies the identity

$$a(\mathbf{u}, \mathbf{v}) = (\mathbf{f}, \mathbf{v})_0 + b(\mathbf{h}, \mathbf{v}) \tag{9}$$

for all test functions  $\mathbf{v} \in V$ , the equation of continuity (3) a.e. in  $\Omega$  and the boundary conditions (3)–(8) (on  $\Gamma_i$ ,  $\Gamma_w$ ,  $\Gamma_-$  and  $\Gamma_+$ ) in the sense of traces.

## 1.4 Existence and uniqueness of a weak solution

**Lemma 1.** *There exists a constant  $c_1 > 0$  independent of  $\mathbf{g}$  and a divergence-free extension  $\mathbf{g}^* \in H^1(\Omega)^2$  of function  $\mathbf{g}$  from  $\Gamma_i$  onto  $\Omega$  such that  $\mathbf{g}^* = \mathbf{0}$  on  $\Gamma_w$ ,  $\mathbf{g}^*$  satisfies the condition of periodicity*

$$\mathbf{g}^*(x_1, x_2 + \tau) = \mathbf{g}^*(x_1, x_2) \quad \text{for } (x_1, x_2) \in \Gamma_- \quad (10)$$

and the estimate

$$\|\mathbf{g}^*\|_1 \leq c_1 \|\mathbf{g}\|_{s; \Gamma_i}. \quad (11)$$

Now we construct the weak solution  $\mathbf{u}$  in the form  $\mathbf{u} = \mathbf{g}^* + \mathbf{z}$  where  $\mathbf{z} \in V$  is a new unknown function. Substituting  $\mathbf{u} = \mathbf{g}^* + \mathbf{z}$  into equation (9), we get the following problem: *Find a function  $\mathbf{z} \in V$  such that it satisfies the equation*

$$a(\mathbf{g}^* + \mathbf{z}, \mathbf{v}) = (\mathbf{f}, \mathbf{v})_0 + b(\mathbf{h}, \mathbf{v}) \quad (12)$$

for all  $\mathbf{v} \in V$ .

**Theorem (on the existence of a weak solution).** *There exists  $\varepsilon > 0$  such that if  $\|\mathbf{g}\|_{s;\Gamma_i} < \varepsilon$  then there exists a solution  $\mathbf{u}$  of the problem defined in Definition 1.*

### Principle of the proof.

We use the Galerkin method and we construct approximations  $\mathbf{z}_n$  in  $n$ -dimensional subspaces  $V_n$  of  $V$ . The fundamental tool which guarantees the existence of the approximations is the **coerciveness** of the bilinear form  $a$  in space  $V$ . Applying successively estimates of the forms  $a_1$ ,  $a_2$  and  $b$ , we can derive the next lemma.

**Lemma 2.** *There exist positive constants  $c_2$  and  $c_3$  such that*

$$a(\mathbf{g}^* + \mathbf{z}, \mathbf{z}) \geq \|\mathbf{z}\| \left( \nu \|\mathbf{z}\| - \nu c_4 c_5 \|\mathbf{g}\|_{H^s(\Gamma_i)^2} - c_2 \|\mathbf{g}\|_{H^s(\Gamma_i)^2}^2 - c_3 \|\mathbf{g}\|_{H^s(\Gamma_i)^2} \|\mathbf{z}\| \right) \quad (13)$$

Now the coerciveness of the form  $a$  follows from (13) and the assumption on a sufficient smallness of  $\|\mathbf{g}\|_{s;\Gamma_i}$ .

**Theorem (on the uniqueness of a weak solution).** *There exists  $R > 0$  such that if  $\mathbf{u}_1$  and  $\mathbf{u}_2$  are two solutions of the problem from Definition 1 such that  $\|\nabla \mathbf{u}_1\|_{L^2(\Omega)^4} \leq R$  then  $\mathbf{u}_1 = \mathbf{u}_2$ .*

**Principle of the proof.** The structure of the form  $a_2$  implies that

$$a_2(\mathbf{u}, \mathbf{v}, \mathbf{v}) = 0, \quad \text{for all } \mathbf{u} \text{ and } \mathbf{v} \text{ from } H^1(\Omega)^2.$$

$$a(\mathbf{u}_1, \mathbf{v}) - a(\mathbf{u}_2, \mathbf{v}) = 0.$$

$$a_1(\mathbf{u}_1 - \mathbf{u}_2, \mathbf{u}_1 - \mathbf{u}_2) = \nu \|\|\mathbf{u}_1 - \mathbf{u}_2\|\|^2$$

$$|a_2(\mathbf{u}_1, \mathbf{u}_1, \mathbf{u}_1 - \mathbf{u}_2) - a_2(\mathbf{u}_2, \mathbf{u}_2, \mathbf{u}_1 - \mathbf{u}_2)| \leq c_6 R \|\|\mathbf{u}_1 - \mathbf{u}_2\|\|^2$$

$$\nu \|\|\mathbf{u}_1 - \mathbf{u}_2\|\|^2 \leq c_6 R \|\|\mathbf{u}_1 - \mathbf{u}_2\|\|^2.$$

## 2. The problem with a nonlinear boundary condition on the outflow

### 2.1. Equations of motion

The conservation of momentum is expressed by the **Navier–Stokes equation** in the form

$$(\mathbf{u} \cdot \nabla)\mathbf{u} = \mathbf{f} - \nabla p + \nu \Delta \mathbf{u}. \quad (14)$$

The condition of incompressibility is expressed by **the equation of continuity**

$$\operatorname{div} \mathbf{u} = 0. \quad (15)$$

## 2.2 Boundary conditions

**Boundary conditions** on  $\Gamma_i, \Gamma_+, \Gamma_-$  and  $\Gamma_w$  are the same as in Section 1.

**The nonlinear condition on the outflow  $\Gamma_o$ :**

$$-\nu \frac{\partial \mathbf{u}}{\partial \mathbf{n}} + p \mathbf{n} - \frac{1}{2} (\mathbf{u} \cdot \mathbf{n})^- \mathbf{u} = \mathbf{h}. \quad (16)$$

## 2.3 The weak formulation

The used integral identity now has the form

$$a(\mathbf{u}, \mathbf{v}) = (\mathbf{f}, \mathbf{v})_{L^2(\Omega)^2} + b(\mathbf{h}, \mathbf{v}) \quad \forall \mathbf{v} \in V,$$

where

$$a(\mathbf{u}, \mathbf{v}) \quad := \quad a_1(\mathbf{u}, \mathbf{v}) + a_2(\mathbf{u}, \mathbf{u}, \mathbf{v}) + a_3(\mathbf{u}, \mathbf{u}, \mathbf{v}),$$

$$a_1(\mathbf{u}, \mathbf{v}) \quad := \quad \nu \int_{\Omega} \sum_{i,j=1}^2 \frac{\partial u_i}{\partial x_j} \frac{\partial v_i}{\partial x_j} \, d\mathbf{x},$$

$$a_2(\mathbf{u}, \mathbf{v}, \mathbf{w}) \quad := \quad \int_{\Omega} \sum_{i,j=1}^2 u_j \frac{\partial v_i}{\partial x_j} w_i \, d\mathbf{x},$$

$$a_3(\mathbf{u}, \mathbf{v}, \mathbf{w}) \quad := \quad \int_{\Gamma_o} \frac{1}{2} (\mathbf{u} \cdot \mathbf{n})^- \mathbf{v} \cdot \mathbf{w} \, dS,$$

$$b(\mathbf{h}, \mathbf{v}) \quad := \quad - \int_{\Gamma_o} \mathbf{h} \cdot \mathbf{v} \, dS.$$

The definition of the weak solution is formally identical with Definition 1, we only use form  $a$  in the form given on the preceding page.

**Definition 2.** Let  $\mathbf{g} \in H^s(\Gamma_i)^2$  (for some  $s \in (\frac{1}{2}, 1]$ ) satisfy the condition  $\mathbf{g}(A_1) = \mathbf{g}(A_0)$ . Let  $\mathbf{f} \in L^2(\Omega)^2$  and  $\mathbf{h} \in L^2(\Gamma_o)^2$ . The **weak solution** of the problem (14), (15), (3)–(7), (16) is a vector function  $\mathbf{u} \in H^1(\Omega)^2$  which satisfies the identity

$$a(\mathbf{u}, \mathbf{v}) = (\mathbf{f}, \mathbf{v})_0 + b(\mathbf{h}, \mathbf{v}) \quad (17)$$

for all test functions  $\mathbf{v} \in V$ , the equation of continuity (15) a.e. in  $\Omega$  and the boundary conditions (3)–(7) (on  $\Gamma_i$ ,  $\Gamma_w$ ,  $\Gamma_-$  and  $\Gamma_+$ ) in the sense of traces.

**Theorem (on the existence of a weak solution).** There exists  $\varepsilon > 0$  such that if  $\|\mathbf{g}\|_{s;\Gamma_i} < \varepsilon$  then there exists a solution  $\mathbf{u}$  of the problem defined in Definition 2.



**Theorem (on the uniqueness of a weak solution).** *There exists  $R > 0$  such that if  $\mathbf{u}_1$  and  $\mathbf{u}_2$  are two solutions of the problem from Definition 2 such that  $\|\nabla \mathbf{u}_1\|_{L^2(\Omega)^4} \leq R$  and  $\|\nabla \mathbf{u}_2\|_{L^2(\Omega)^4} \leq R$  then  $\mathbf{u}_1 = \mathbf{u}_2$ .*

**Principle of the proof.**

$$a(\mathbf{u}_1, \mathbf{v}) - a(\mathbf{u}_2, \mathbf{v}) = 0.$$

$$a_1(\mathbf{u}_1 - \mathbf{u}_2, \mathbf{u}_1 - \mathbf{u}_2) = \nu \|\|\mathbf{u}_1 - \mathbf{u}_2\|\|^2$$

$$|a_2(\mathbf{u}_1, \mathbf{u}_1, \mathbf{u}_1 - \mathbf{u}_2) - a_2(\mathbf{u}_2, \mathbf{u}_2, \mathbf{u}_1 - \mathbf{u}_2)| \leq c_7 R \|\|\mathbf{u}_1 - \mathbf{u}_2\|\|^2$$

$$|a_3(\mathbf{u}_1, \mathbf{u}_1, \mathbf{u}_1 - \mathbf{u}_2) - a_3(\mathbf{u}_2, \mathbf{u}_2, \mathbf{u}_1 - \mathbf{u}_2)| \leq c_8 R \|\|\mathbf{u}_1 - \mathbf{u}_2\|\|^2$$

$$\nu \|\|\mathbf{u}_1 - \mathbf{u}_2\|\|^2 \leq c_7 + c_8 R \|\|\mathbf{u}_1 - \mathbf{u}_2\|\|^2.$$

The derivation of the estimate for  $a_3$  is technically more complicated.